

A NEW PROOF OF MOSER'S THEOREM

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Abstract In this paper, we consider the persistence of invariant tori for mappings under perturbations. Mainly, we give a new proof of Moser type theorem about invariant tori for twist mappings with intersection property.

Keywords High-dimensional Moser's theorem, twist mapping, intersection property.

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1. Introduction

In 1954, Kolmogorov [7] first claimed the persistence of quasi-periodic invariant tori under small perturbations with non-degenerate conditions and Arnold [1] proved it soon after. In 1962, Moser [11] proved the existence of the invariant curves for twist mappings with one action and one angular variable. These works [1, 7, 11] established the celebrated KAM theory. Later, Zehnder [22, 23] took use of the generalized implicit function theorem to solve some divisor problems. Rüssmann [16] exhibited the optimal result that the invariant tori of twist mappings persist when the perturbations belong to class C^l , $l > 3$. Due to the efforts of Cheng and Sun [2], the persistence of invariant tori for three-dimensional measure-preserving mappings under small perturbations is proved. And the case of volume-preserving mappings on n -dimensional space was considered by Xia [20]. For Moser's invariant curve theorem, one can find some further developments in Svanidze [17], Herman [6], Levi [9], and Moser [12–14].

Cong et al. [3] proved the persistence of invariant tori of analytic nearly twist mappings with different numbers of action and angular variables when having intersection property. The so-called intersection property is that any n -dimensional torus close to the invariant torus of the unperturbed system intersects its image under the mapping. Li and Yi [10] introduced a parameter in their consideration of generalized Hamiltonian systems. It inspired us to consider the corresponding circumstances of twist mappings with a parameter. Recently, Yang and Li [21] investigated the existence of periodically invariant tori of twist mappings on resonant surfaces. There are more related works on KAM theory in recent years, for example, see Chierchia et al. [4], Calleja et al. [5], Koudjina [8], Sevryuk [18], and Trujillo [19]. In this paper, we will show detailed proof of the persistence of invariant tori for twist mappings with intersection property.

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Consider the analytic mapping $\mathcal{F} : T^n \times D \times G \rightarrow T^n \times \mathbb{R}^{m+p}$ defined by

$$\mathcal{F} : \begin{cases} \theta^1 = \theta + \omega(r, \xi) + f(\theta, r, \xi), \\ r^1 = r + g(\theta, r, \xi), \end{cases} \quad (1.1)$$

where $\theta \in T^n = \mathbb{R}^n / 2\pi\mathbb{Z}^n$, $r \in D \subset \mathbb{R}^m$, $\xi \in G \subset \mathbb{R}^p$, D and G are bounded and closed domains. Assume the following conditions hold:

- (A1) f and g are analytic on $T^n \times D \times G$;
- (A2) there exists a positive integer K such that

$$\text{rank} \left\{ \frac{\partial^\alpha \omega}{\partial(r, \xi)^\alpha} : 0 \leq |\alpha| \leq K-1 \right\} = n;$$

- (A3) the frequency ω satisfies the Diophantine condition

$$|\langle k, \omega \rangle - k_0| \geq \frac{\delta}{|k|^\tau}, \quad \forall k \in \mathbb{Z}^n \setminus \{0\}, \quad \forall k_0 \in \mathbb{Z},$$

where $\tau > n+1$ is fixed.

The unperturbed mapping of (1.1) is

$$\begin{cases} \theta^1 = \theta + \omega(r, \xi), \\ r^1 = r, \end{cases} \quad (1.2)$$

which is called twist mapping.

Hereafter consider the mapping $\mathcal{F}$ in a complex neighborhood $\mathbb{T}^n \times \mathbb{D} \times \mathbb{G}$ of $T^n \times D \times G$, where $\mathbb{T}^n = \mathbb{C}^n / 2\pi\mathbb{Z}^n$, $\mathbb{D} \subset \mathbb{C}^m$, and $\mathbb{G} \subset \mathbb{C}^p$. For given $0 < \rho, s, \sigma \leq 1$, let

$$\mathcal{D}(\rho, s, \sigma) = \{(\theta, r, \xi) \in \mathbb{T}^n \times \mathbb{D} \times \mathbb{G} : |\text{Im } \theta| \leq \rho, |\text{Im } r| \leq s, |\text{Im } \xi| \leq \sigma\},$$

and $\|\cdot\|_{\mathcal{D}(\rho, s, \sigma)}$ represents the sup-norm of functions in $\mathcal{D}(\rho, s, \sigma)$.

Now we are in a position to state the main results of this paper.

Theorem 1.1. *Consider the twist mapping (1.1) with intersection property on $T^n \times D \times G$. Assume (A1)-(A3) hold. Then there exists a nonempty Cantor set $D_\delta \subset D \times G$ and a sufficiently small constant $\epsilon_* > 0$ such that for all $\epsilon \in (0, \epsilon_*)$, if*

$$\|f(\theta, r, \xi)\|_{\mathcal{D}(\rho, s, \sigma)} + \|g(\theta, r, \xi)\|_{\mathcal{D}(\rho, s, \sigma)} \leq \epsilon, \quad (1.3)$$

the mapping $\mathcal{F}$ has a Whitney smooth family of invariant tori for $(r, \xi) \in D_\delta \subset D \times G$. And the measure estimate holds:

$$\text{meas}((D \times G) \setminus D_\delta) = \mathcal{O}(\delta^{\frac{1}{K}}).$$

The case $m = 0$ in Theorem 1.1 can be regarded as a parameterized Herman's theorem under Rüssmann type nondegenerate condition. See the corollary as follows.

Corollary 1.1. *Consider the analytic mapping $\mathcal{F}'$ defined by*

$$\theta^1 = \theta + \omega(\xi) + f(\theta, \xi), \quad (1.4)$$

where $\theta \in T^n = \mathbb{R}^n/2\pi\mathbb{Z}^n$, $\xi \in G \subset \mathbb{R}^p$ is a parameter and G is a bounded and closed domain. Assume (A1), (A3) and (A2)' hold, where (A2)' there exists a positive integer K such that

$$\text{rank}\left\{\frac{\partial^\alpha \omega}{\partial \xi^\alpha} : 0 \leq |\alpha| \leq K-1\right\} = n.$$

Then there exist a complex domain $\mathcal{D}(\rho, \sigma) = \{(\theta, \xi) \in \mathbb{T}^n \times \mathbb{G} : |\text{Im } \theta| \leq \rho, |\text{Im } \xi| \leq \sigma\}$, a nonempty Cantor set $G_\delta \subset G$, and a sufficiently small constant $\epsilon_* > 0$ such that for all $\epsilon \in (0, \epsilon_*)$, if

$$\|f(\theta, \xi)\|_{\mathcal{D}(\rho, \sigma)} \leq \epsilon,$$

then mapping $\mathcal{F}'$ has a Whitney smooth family of invariant tori. And the measure estimate holds:

$$\text{meas}(G \setminus G_\delta) = \mathcal{O}(\delta^{\frac{1}{K}}).$$

The rest of this paper is organized as follows. Section 2 gives general analytic KAM steps for twist mappings with intersection property; section 3 draws up the proof of Theorem 1.1. In Appendix, we introduce the method of measure estimates mentioned in [15].

2. KAM Steps

In this section, mainly, we consider the KAM steps for mapping (1.1). The symbol $'$ represents the abbreviation of constants appearing in estimates for convenience that are independent of the iterative process.

2.1. Description of the 0-th KAM step

We first define the following parameters that appear in the 0-th KAM step.

$$\begin{aligned} s_0 &= \epsilon_0^{\frac{1}{2K}}, & \delta_0 &= \epsilon_0^{\frac{a_0}{4K}}, & \mu_0 &= \epsilon_0^{\frac{1}{2} - \frac{a_0}{4}}, \\ \epsilon_0 &= \epsilon = s_0^K \delta_0^K \mu_0, & \rho_0 &= \rho, & \sigma_0 &= \epsilon_0^{\frac{1}{2K}}, \\ \mathcal{D}(\rho_0, s_0, \sigma_0) &= \{(\theta_0, r_0, \xi) : |\text{Im } \theta_0| \leq \rho_0, |\text{Im } r_0| \leq s_0, |\text{Im } \xi| \leq \sigma_0\}, \end{aligned}$$

where $0 < a_0 < 2$. We also define $\rho_\nu = (\frac{1}{2} + \frac{1}{2^{\nu+1}})\rho_0$, $s_\nu = \frac{1}{2^{\nu+2}}s_0$, $\sigma_\nu = \frac{1}{2^{\nu+2}}\sigma_0$, ($\nu = 1, 2, \dots$). And χ is an integer that satisfies $2 < (\frac{4}{3})^\chi < 4$, for instance, $\chi = 3$ or 4.

The mapping at 0-th step is

$$\mathcal{F}_0 : \begin{cases} \theta_0^1 = \theta_0 + \omega_0(r_0, \xi) + f_0, \\ r_0^1 = r_0 + g_0, \end{cases} \quad (2.1)$$

where $(\theta_0, r_0, \xi) \in \mathcal{D}(\rho_0, s_0, \sigma_0)$, $\omega_0 = \omega$ is the original frequency vector, and $f_0 = f$, $g_0 = g$ are original perturbations.

From (1.3) we have

$$\|f_0(\theta_1, r_1, \xi)\|_{\mathcal{D}(\rho_0, s_0, \sigma_0)} + \|g_0(\theta_1, r_1, \xi)\|_{\mathcal{D}(\rho_0, s_0, \sigma_0)} \leq s_0^K \delta_0^K \mu_0.$$

For $0 \leq \kappa, j, h \leq K$, $\kappa + j + h = K$, denote $\partial^K(\cdot) = \partial_\theta^\kappa \partial_r^j \partial_\xi^h(\cdot)$ throughout this paper. Then taking use of Cauchy estimate on $\hat{\mathcal{D}}(\rho_1, s_1, \sigma_1) = \mathcal{D}(\rho_1 + \frac{3}{4}(\rho_0 - \rho_1), s_1, \sigma_1)$, we get

$$\begin{aligned} & \| \partial^K f_0(\theta_1, r_1, \xi) \|_{\hat{\mathcal{D}}(\rho_1, s_1, \sigma_1)} + \| \partial^K g_0(\theta_1, r_1, \xi) \|_{\hat{\mathcal{D}}(\rho_1, s_1, \sigma_1)} \\ &= \| \partial_{\theta_1}^\kappa \partial_{r_1}^j \partial_\xi^h f_0(\theta_1, r_1, \xi) \|_{\hat{\mathcal{D}}(\rho_1, s_1, \sigma_1)} + \| \partial_{\theta_1}^\kappa \partial_{r_1}^j \partial_\xi^h g_0(\theta_1, r_1, \xi) \|_{\hat{\mathcal{D}}(\rho_1, s_1, \sigma_1)} \\ &\leq \frac{4^\kappa \kappa! j! h! (\| f_0(\theta_1, r_1, \xi) \|_{\mathcal{D}(\rho_0, s_0, \sigma_0)} + \| g_0(\theta_1, r_1, \xi) \|_{\mathcal{D}(\rho_0, s_0, \sigma_0)})}{(\rho_0 - \rho_1)^\kappa (s_0 - s_1)^j (\sigma_0 - \sigma_1)^h} \\ &\leq \frac{s_0^{K-j} \delta_0^K \mu_0}{(\rho_0 - \rho_1)^\kappa (\sigma_0 - \sigma_1)^h}. \end{aligned}$$

2.2. Description of the ν -th KAM step

In this subsection, we will show one cycle of KAM steps from ν -th step to $\nu + 1$. For any $\nu = 1, 2, \dots$, set

$$\begin{aligned} \epsilon_\nu &= s_\nu^K \delta_\nu^K \mu_\nu, \quad s_\nu = \frac{1}{2^{7\nu+2}} s_0, \quad \delta_\nu = \left(\frac{1}{2} + \frac{1}{2^{\nu+1}}\right) \delta_0, \quad \mu_\nu = \mu_{\nu-1}^{\frac{4}{3}}, \\ \rho_\nu &= \left(\frac{1}{2} + \frac{1}{2^{\nu+1}}\right) \rho_0, \quad \sigma_\nu = \frac{1}{2^{7\nu+2}} \sigma_0, \quad N_{\nu+1} = ([\log \frac{1}{\mu_\nu}] + 1)^{3\chi}, \\ \mathcal{D}_j &= \mathcal{D}(\rho_{\nu+1} + \frac{j-1}{4}(\rho_\nu - \rho_{\nu+1}), j s_{\nu+1}, j \sigma_{\nu+1}), \quad j = 1, 2, 3, 4, \\ \hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1}) &= \mathcal{D}(\rho_{\nu+1} + \frac{3}{4}(\rho_\nu - \rho_{\nu+1}), s_{\nu+1}, \sigma_{\nu+1}), \\ \mathcal{D}_1 &= \mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1}). \end{aligned}$$

Suppose that after ν KAM steps, we have arrived at the following mapping

$$\mathcal{F}_\nu : \begin{cases} \theta_\nu^1 = \theta_\nu + \omega_\nu(r_\nu, \xi) + f_\nu(\theta_\nu, r_\nu, \xi), \\ r_\nu^1 = r_\nu + g_\nu(\theta_\nu, r_\nu, \xi), \end{cases} \quad (2.2)$$

where $(\theta_\nu, r_\nu, \xi) \in \mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu) = \{(\theta_\nu, r_\nu, \xi) : |\operatorname{Im} \theta_\nu| \leq \rho_\nu, |\operatorname{Im} r_\nu| \leq s_\nu, |\operatorname{Im} \xi| \leq \sigma_\nu\}$. Introduce a transformation $\mathcal{U}_{\nu+1}$:

$$\begin{cases} \theta_\nu = \theta_{\nu+1} + U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi), \\ r_\nu = r_{\nu+1} + V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi), \end{cases} \quad (2.3)$$

which satisfies

$$\mathcal{U}_{\nu+1} \circ \mathcal{F}_{\nu+1} = \mathcal{F}_\nu \circ \mathcal{U}_{\nu+1}. \quad (2.4)$$

From (2.4) we have

$$\begin{aligned} & \theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi) + f_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \\ & + U_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi) + f_{\nu+1}, r_{\nu+1} + g_{\nu+1}, \xi) \\ &= \theta_\nu + \omega_\nu(r_\nu, \xi) + f_\nu(\theta_\nu, r_\nu, \xi) \\ &= \theta_{\nu+1} + U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + \omega_\nu(r_\nu, \xi) + f_\nu(\theta_{\nu+1} + U_{\nu+1}, r_{\nu+1} + V_{\nu+1}, \xi), \\ & \quad r_{\nu+1} + g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + V_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi) + f_{\nu+1}, r_{\nu+1} + g_{\nu+1}, \xi) \\ &= r_\nu + g_\nu(\theta_\nu, r_\nu, \xi) \end{aligned}$$

$$= r_{\nu+1} + V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + g_{\nu}(\theta_{\nu+1} + U_{\nu+1}, r_{\nu+1} + V_{\nu+1}, \xi),$$

i.e.

$$\begin{aligned} & \omega_{\nu+1}(r_{\nu+1}, \xi) + f_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \\ & + U_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1} + f_{\nu+1}, r_{\nu+1} + g_{\nu+1}, \xi) \\ & = U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + \omega_{\nu}(r_{\nu}, \xi) + f_{\nu}(\theta_{\nu+1} + U_{\nu+1}, r_{\nu+1} + V_{\nu+1}, \xi), \end{aligned} \quad (2.5)$$

$$\begin{aligned} & g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + V_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1} + f_{\nu+1}, r_{\nu+1} + g_{\nu+1}, \xi) \\ & = V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + g_{\nu}(\theta_{\nu+1} + U_{\nu+1}, r_{\nu+1} + V_{\nu+1}, \xi). \end{aligned} \quad (2.6)$$

2.2.1. Estimates of the remainders

Consider the Fourier series expansion of $f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi)$,

$$f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \sum_{k \in \mathbb{Z}^n} f_{k,\nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}. \quad (2.7)$$

The operators of truncation and remainder are respectively

$$\Gamma_{N_{\nu+1}} f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \sum_{0 < |k| \leq N_{\nu+1}} f_{k,\nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}, \quad (2.8)$$

$$R_{N_{\nu+1}} f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \sum_{|k| > N_{\nu+1}} f_{k,\nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}, \quad (2.9)$$

where $f_{k,\nu}(r_{\nu+1}, \xi) = \int_{\mathbb{T}^n} f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) e^{-i\langle k, \theta_{\nu+1} \rangle} d\theta_{\nu+1}$ is the Fourier coefficient of $f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi)$.

Lemma 2.1. *If*

$$(H1) \quad \int_{N_{\nu+1}}^{+\infty} l^n e^{-l \frac{\rho_{\nu} - \rho_{\nu+1}}{4}} dl \leq \mu_{\nu}, \quad (2.10)$$

then

$$\| R_{N_{\nu+1}} f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_3} \leq s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2. \quad (2.11)$$

Proof. By a direct computation, we have

$$\begin{aligned} & \| R_{N_{\nu+1}} f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_3} \\ & \leq \left| \sum_{|k| > N_{\nu+1}} f_{k,\nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle} \right| \\ & \leq \sum_{|k| > N_{\nu+1}} |f_{k,\nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}| \\ & \leq \sum_{|k| > N_{\nu+1}} |f_{k,\nu}(r_{\nu+1}, \xi)| |e^{i\langle k, \theta_{\nu+1} \rangle}| \\ & \leq \| f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4} \sum_{|k| > N_{\nu+1}} e^{-|k|(\rho_{\nu+1} + \frac{3}{4}(\rho_{\nu} - \rho_{\nu+1}))} e^{|k|(\rho_{\nu+1} + \frac{1}{2}(\rho_{\nu} - \rho_{\nu+1}))} \\ & \leq \| f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4} \sum_{|k| > N_{\nu+1}} e^{-|k| \frac{\rho_{\nu} - \rho_{\nu+1}}{4}} \end{aligned}$$

$$\begin{aligned}
&\leq \|f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}_4} \sum_{|k| > N_{\nu+1}} |k|^n e^{-|k| \frac{\rho_\nu - \rho_{\nu+1}}{4}} \\
&\leq \|f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}_4} \int_{N_{\nu+1}}^{+\infty} l^n e^{-l \frac{\rho_\nu - \rho_{\nu+1}}{4}} dl \\
&\leq s_\nu^K \delta_\nu^K \mu_\nu^2.
\end{aligned}$$

□

Remark 2.1. On the third inequality, we notice that the coefficients $f_{k,\nu}(r_{\nu+1}, \xi)$ decay exponentially, that is,

$$|f_{k,\nu}(r_{\nu+1}, \xi)| \leq \|f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}_4} e^{-|k|(\rho_{\nu+1} + \frac{3}{4}(\rho_\nu - \rho_{\nu+1}))}. \quad (2.12)$$

From Cauchy estimate we get

$$\begin{aligned}
&\|\partial^K R_{N_{\nu+1}} f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&= \|\partial_{\theta_{\nu+1}}^\kappa \partial_{r_{\nu+1}}^j \partial_\xi^h R_{N_{\nu+1}} f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&\leq \cdot \frac{\|R_{N_{\nu+1}} f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu)}}{(\frac{\rho_\nu - \rho_{\nu+1}}{4})^\kappa (s_\nu - s_{\nu+1})^j (\sigma_\nu - \sigma_{\nu+1})^h} \\
&\leq \cdot \frac{s_\nu^K \delta_\nu^K \mu_\nu^2}{(\rho_\nu - \rho_{\nu+1})^\kappa (\sigma_\nu - \sigma_{\nu+1})^h s_{\nu+1}^j} \\
&\leq \cdot \frac{s_\nu^{K-j} \delta_\nu^K \mu_\nu^2}{(\rho_\nu - \rho_{\nu+1})^\kappa (\sigma_\nu - \sigma_{\nu+1})^h}. \quad (2.13)
\end{aligned}$$

Remark 2.2. Note

$$\int_{N_{\nu+1}}^{+\infty} l^n e^{-l \frac{\rho_\nu - \rho_{\nu+1}}{4}} dl = \sum_{i=0}^n \frac{4^{i+1} \overbrace{n(n-1) \cdots (n-i+1)}^i}{(\rho_\nu - \rho_{\nu+1})^{i+1}} N_{\nu+1}^{n-i} e^{-N_{\nu+1} \frac{\rho_\nu - \rho_{\nu+1}}{4}}.$$

Then in order to get

$$\int_{N_{\nu+1}}^{+\infty} l^n e^{-l \frac{\rho_\nu - \rho_{\nu+1}}{4}} dl < \mu_\nu,$$

it suffices to hold

$$\frac{4^{n+1} (n+1)!}{(\rho_\nu - \rho_{\nu+1})^{n+1}} N_{\nu+1}^n e^{-N_{\nu+1} \frac{\rho_\nu - \rho_{\nu+1}}{4}} < \mu_\nu,$$

i.e.

$$\begin{aligned}
&(n+1)\log 4 + \log(n+1)! + (\nu+2)(n+1)\log 2 - (n+1)\log \rho_0 \\
&+ 3\chi n \log([\log \frac{1}{\mu_\nu}] + 1) - \frac{1}{2^{\nu+4}} \rho_0 ([\log \frac{1}{\mu_\nu}] + 1)^{3\chi} \\
&< \log \mu_\nu.
\end{aligned}$$

Since $2 < (\frac{4}{3})^\chi < 4$, we get

$$\frac{\rho_0}{2^{\nu+4}} ([\log \frac{1}{\mu_\nu}] + 1)^\chi \geq \frac{\rho_0}{2^{\nu+4}} \left((\frac{4}{3})^\nu [\log \frac{1}{\mu_0}] \right)^\chi$$

$$\begin{aligned}
&\geq \frac{\rho_0}{2^{\nu+4}} \left(\frac{4}{3}\right)^{\chi\nu} [\log \frac{1}{\mu_0}]^\chi \\
&\geq -\frac{\rho_0}{2^{\nu+4}} \left(\frac{4}{3}\right)^{\chi\nu} [\log \mu_0]^\chi \\
&\geq -\frac{\rho_0}{2^4} [\log \mu_0]^\chi \\
&\geq 1.
\end{aligned}$$

It follows that

$$\begin{aligned}
&(n+1)\log 4 + \log(n+1)! - (n+1)\log \rho_0 + (n+1)(\nu+2)\log 2 \\
&\quad + 3n\chi \log([\log \frac{1}{\mu_\nu}] + 1) - \frac{\rho_0}{2^{\nu+4}} ([\log \frac{1}{\mu_\nu}] + 1)^{3\chi} \\
&\leq (n+1)\log 4 + \log(n+1)! - (n+1)\log \rho_0 + (n+1)(\nu+2)\log 2 \\
&\quad + 3n\chi \log([\log \frac{1}{\mu_\nu}] + 1) - [\log \frac{1}{\mu_\nu}]^{2\chi} \\
&\leq -\log \frac{1}{\mu_\nu}.
\end{aligned}$$

2.2.2. Homological equations

Before considering the homological equations in KAM steps, it is necessary to take time to consider the Diophantine condition in ν -th step to make sure that $\mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu)$ is well defined.

Lemma 2.2. *There exist positive constants M_0 , M_1 , and M_2 such that $|\omega_\nu|$, $|\frac{\partial \omega_\nu}{\partial r_\nu}|$, and $|\frac{\partial \omega_\nu}{\partial \xi}|$ are bounded by them on $\mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu)$ respectively. If*

$$(H2) \quad (M_1 + M_2)([\log \frac{1}{\mu_0}] + 1)^{3\chi} \delta_0^{\frac{2}{a_0} - 1} \leq \frac{1}{|k|^\tau},$$

the frequency ω_ν satisfies

$$|\langle k, \omega_\nu(r_\nu, \xi) \rangle - k_0| \geq \frac{\delta_\nu}{2|k|^\tau}, \quad \forall |k| \in (0, N_{\nu+1}], \quad \forall k_0 \in \mathbb{Z}. \quad (2.14)$$

Proof. Due to

$$\omega_\nu(r_\nu, \xi) = \omega_0(r_0, \xi) + \sum_{l=0}^{\nu-1} \bar{f}_l(r_{l+1}, \xi), \quad (2.15)$$

the $\bar{f}_l(r_{l+1}, \xi)$ above is the average of $f_l(\theta_{l+1}, r_{l+1}, \xi)$ and satisfies

$$\begin{aligned}
|\bar{f}_l(r_{l+1}, \xi)| &= \frac{1}{(2\pi)^n} \left| \int_{\mathbb{T}^n} f_l(\theta_{l+1}, r_{l+1}, \xi) d\theta_{l+1} \right| \\
&\leq \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n} \|f_l(\theta_{l+1}, r_{l+1}, \xi)\|_{\mathcal{D}(\rho_l, s_l, \sigma_l)} d\theta_{l+1} \\
&\leq s_l^K \delta_l^K \mu_l,
\end{aligned}$$

and then ω_ν is well defined in its domain with

$$|\omega_\nu| \leq |\omega_0| + \sum_{l=0}^{\nu-1} |\bar{f}_l(\theta_{l+1}, r_{l+1}, \xi)| := M_0.$$

Therefore, the bounds of $|\frac{\partial \omega_\nu}{\partial r_\nu}|$ and $|\frac{\partial \omega_\nu}{\partial \xi}|$ are respectively,

$$\begin{aligned} |\frac{\partial \omega_\nu}{\partial r_\nu}| &\leq |\frac{\partial \omega_0}{\partial r_\nu}| + \sum_{l=0}^{\nu-1} |\frac{\partial \bar{f}_l(r_l, \xi)}{\partial r_\nu}| \\ &\leq |\frac{\partial \omega_0}{\partial r_\nu}| + \sum_{l=0}^{\nu-1} \frac{s_l^K \delta_l^K \mu_l}{s_l - s_{l+1}} \\ &:= M_1, \\ |\frac{\partial \omega_\nu}{\partial \xi}| &\leq |\frac{\partial \omega_0}{\partial \xi}| + \sum_{l=0}^{\nu-1} |\frac{\partial \bar{f}_l(r_l, \xi)}{\partial \xi}| \\ &\leq |\frac{\partial \omega_0}{\partial \xi}| + \sum_{l=0}^{\nu-1} \frac{s_l^K \delta_l^K \mu_l}{\sigma_l - \sigma_{l+1}} \\ &:= M_2. \end{aligned}$$

We see that

$$\begin{aligned} N_{\nu+1} &= ([\log \frac{1}{\mu_\nu}] + 1)^{3\chi} \\ &= \left(\left(\frac{4}{3} \right)^\nu [\log \frac{1}{\mu_0}] + 1 \right)^{3\chi} \\ &\leq \left(\frac{4}{3} \right)^{3\chi\nu} ([\log \frac{1}{\mu_0}] + 1)^{3\chi} \\ &\leq 2^{6\nu} ([\log \frac{1}{\mu_0}] + 1)^{3\chi}. \end{aligned}$$

Moreover, for $(\theta_{\nu+1}, r_{\nu+1}, \xi)$ and $(\theta'_{\nu+1}, r'_{\nu+1}, \xi')$ belonging to $\mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu)$, we have

$$\begin{aligned} &|\langle k, \omega_\nu(r_\nu, \xi) - \omega_\nu(r'_\nu, \xi') \rangle| \\ &\leq |\langle k, \omega_\nu(r_\nu, \xi) - \omega_\nu(r'_\nu, \xi) + \omega_\nu(r'_\nu, \xi) - \omega_\nu(r'_\nu, \xi') \rangle| \\ &\leq |\langle k, \omega_\nu(r_\nu, \xi) - \omega_\nu(r'_\nu, \xi) \rangle| + |\langle k, \omega_\nu(r'_\nu, \xi) - \omega_\nu(r'_\nu, \xi') \rangle| \\ &\leq |k| |\omega_\nu(r_\nu, \xi) - \omega_\nu(r'_\nu, \xi)| + |k| |\omega_\nu(r'_\nu, \xi) - \omega_\nu(r'_\nu, \xi')| \\ &\leq |k| |\frac{\partial \omega_\nu}{\partial r_\nu}| |r_\nu - r'_\nu| + |k| |\frac{\partial \omega_\nu}{\partial \xi}| |\xi - \xi'| \\ &\leq M_1 s_\nu |k| + M_2 \sigma_\nu |k| \\ &\leq N_{\nu+1} (M_1 s_\nu + M_2 \sigma_\nu) \\ &\leq 2^{6\nu} ([\log \frac{1}{\mu_0}] + 1)^{3\chi} \frac{1}{2^{7\nu+2}} (M_1 s_0 + M_2 \sigma_0) \\ &= \frac{1}{2^{\nu+2}} ([\log \frac{1}{\mu_0}] + 1)^{3\chi} (M_1 \epsilon_0^{\frac{a_0}{4K} \frac{2}{a_0}} + M_2 \epsilon_0^{\frac{a_0}{4K} \frac{2}{a_0}}) \\ &= \frac{\delta_0}{2^{\nu+2}} ([\log \frac{1}{\mu_0}] + 1)^{3\chi} (M_1 + M_2) \delta_0^{\frac{2}{a_0} - 1} \\ &\leq (\frac{1}{2^2} + \frac{1}{2^{\nu+2}}) \delta_0 ([\log \frac{1}{\mu_0}] + 1)^{3\chi} (M_1 + M_2) \delta_0^{\frac{2}{a_0} - 1} \\ &\leq \frac{\delta_\nu}{2|k|^\tau}. \end{aligned}$$

From the Diophantine condition in ν -th step, one has

$$\begin{aligned}
 |\langle k, \omega_\nu(r_\nu, \xi) \rangle - k_0| &= |\langle k, \omega_\nu(r'_\nu, \xi') + \omega_\nu(r_\nu, \xi) - \omega_\nu(r'_\nu, \xi') \rangle - k_0| \\
 &\geq |\langle k, \omega_\nu(r'_\nu, \xi') \rangle - k_0| - |\langle k, \omega_\nu(r_\nu, \xi) - \omega_\nu(r'_\nu, \xi') \rangle| \\
 &\geq \frac{\delta_\nu}{|k|^\tau} - \frac{\delta_\nu}{2|k|^\tau} \\
 &\geq \frac{\delta_\nu}{2|k|^\tau}.
 \end{aligned}$$

□

From (2.5), (2.6), and (2.15) we now get homological equations

$$\begin{cases} U_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}, r_{\nu+1}, \xi) - U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \Gamma_{N_{\nu+1}} f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi), \\ V_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}, r_{\nu+1}, \xi) - V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \Gamma_{N_{\nu+1}} g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi). \end{cases} \quad (2.16)$$

Then the following lemma yields the solutions of (2.16) and their estimates.

Lemma 2.3. *Assume that*

$$(H3) \quad s_\nu^K \delta_\nu^{K-1} \mu_\nu \leq \frac{1}{4N_{\nu+1}^{\tau+1}}.$$

Then the second equation of (2.16) has a unique analytic solution $V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)$ of zero average on $\mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu)$ and the estimate

$$\|V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}_3} \leq \frac{\|g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}_4}}{\delta_\nu(\rho_\nu - \rho_{\nu+1})^{\tau+n+1}} \quad (2.17)$$

holds.

Proof. The Fourier expansion of $V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)$ is

$$V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \sum_{k \in \mathbb{Z}^n} V_{k, \nu+1}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle},$$

and the truncation of $V_{\nu+1}$ is still represented by $V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)$, that is,

$$V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \sum_{0 < |k| \leq N_{\nu+1}} V_{k, \nu+1}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}.$$

Let the Fourier expansion, truncation, and remainder of $g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)$ be respectively

$$\begin{aligned} g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi) &= \sum_{k \in \mathbb{Z}^n} g_{k, \nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}, \\ \Gamma_{N_{\nu+1}} g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi) &= \sum_{0 < |k| \leq N_{\nu+1}} g_{k, \nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}, \\ R_{N_{\nu+1}} g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi) &= \sum_{|k| > N_{\nu+1}} g_{k, \nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}, \end{aligned}$$

where $g_{k, \nu}(r_{\nu+1}, \xi) = \int_{\mathbb{T}^n} g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi) e^{-i\langle k, \theta_{\nu+1} \rangle} d\theta_{\nu+1}$ is the Fourier coefficient of $g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi)$.

Putting them into the second equation of (2.16), we get

$$\sum_{0 < |k| \leq N_{\nu+1}} V_{k,\nu+1} e^{i\langle k, \theta_{\nu+1} + \omega_{\nu+1} \rangle} - \sum_{0 < |k| \leq N_{\nu+1}} V_{k,\nu+1} e^{i\langle k, \theta_{\nu+1} \rangle} = \sum_{0 < |k| \leq N_{\nu+1}} g_{k,\nu} e^{i\langle k, \theta_{\nu+1} \rangle},$$

i.e.

$$\sum_{0 < |k| \leq N_{\nu+1}} V_{k,\nu+1} e^{i\langle k, \theta_{\nu+1} \rangle} (e^{i\langle k, \omega_{\nu+1} \rangle} - 1) = \sum_{0 < |k| \leq N_{\nu+1}} g_{k,\nu} e^{i\langle k, \theta_{\nu+1} \rangle},$$

that is

$$V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) = \sum_{0 < |k| \leq N_{\nu+1}} \frac{g_{k,\nu}(r_{\nu+1}, \xi) e^{i\langle k, \theta_{\nu+1} \rangle}}{e^{i\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle} - 1}.$$

Then

$$\begin{aligned} & \| V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_3} \\ & \leq \sum_{0 < |k| \leq N_{\nu+1}} \frac{|g_{k,\nu}(r_{\nu+1}, \xi)|}{|e^{i\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle} - 1|} |e^{i\langle k, \theta_{\nu+1} \rangle}| \\ & \leq \sum_{0 < |k| \leq N_{\nu+1}} \frac{|g_{k,\nu}(r_{\nu+1}, \xi)|}{|e^{i\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle} - 1|} e^{|k|(\rho_{\nu+1} + \frac{1}{2}(\rho_{\nu} - \rho_{\nu+1}))} \\ & \leq \frac{2\pi}{\delta_{\nu}} \| g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4} \sum_{0 < |k| \leq N_{\nu+1}} |k|^{\tau} e^{-|k| \frac{\rho_{\nu} - \rho_{\nu+1}}{4}} \\ & \leq \frac{2\pi}{\delta_{\nu}} \| g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4} \int_0^{N_{\nu+1}} l^{\tau+n} e^{-l \frac{\rho_{\nu} - \rho_{\nu+1}}{4}} dl \\ & \leq \frac{2\pi}{\delta_{\nu}} \| g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4} \int_0^{\infty} \left(\frac{4}{\rho_{\nu} - \rho_{\nu+1}} \right)^{\tau+n} x^{\tau+n} e^{-x} d\left(\frac{4x}{\rho_{\nu} - \rho_{\nu+1}} \right) \\ & \leq \frac{2\pi}{\delta_{\nu}} \| g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4} \frac{4^{\tau+n+1}}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} \int_0^{\infty} x^{\tau+n} e^{-x} dx \\ & = \frac{2\pi}{\delta_{\nu}} \| g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4} \frac{4^{\tau+n+1}}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} \Gamma(\tau + n + 1) \\ & \leq \frac{4^{\tau+n+2} \pi(\tau + n)!}{\delta_{\nu}(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} \| g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4}. \end{aligned}$$

□

Remark 2.3. In order to estimate $|e^{i\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle} - 1|$, we first consider $|xe^{i\phi} - y|$ for $x, y \in \mathbb{N}$ and $\phi \in \mathbb{R}$.

$$\begin{aligned} |xe^{i\phi} - y|^2 &= |x \cos \phi + ix \sin \phi - y|^2 \\ &= |x \cos \phi - y|^2 + x^2 \sin^2 \phi \\ &= x^2 \cos^2 \phi + y^2 - 2xy \cos \phi + x^2 \sin^2 \phi \\ &= x^2 + y^2 - 2xy \cos \phi \\ &= x^2 \cos^2 \frac{\phi}{2} + x^2 \sin^2 \frac{\phi}{2} + y^2 \cos^2 \frac{\phi}{2} + y^2 \sin^2 \frac{\phi}{2} \\ &\quad - 2xy \cos^2 \frac{\phi}{2} + 2xy \sin^2 \frac{\phi}{2} \\ &= |x - y|^2 \cos^2 \frac{\phi}{2} + |x + y|^2 \sin^2 \frac{\phi}{2} \end{aligned}$$

$$\begin{aligned} &\geq |x+y|^2 \sin^2 \frac{\phi}{2} \\ &= |x+y|^2 \sin^2 \frac{\phi - 2\pi l}{2}, \quad (l \in \mathbb{Z}). \end{aligned}$$

Therefore,

$$|xe^{i\phi} - y| \geq |x+y| \sin \frac{\phi - 2\pi l}{2}.$$

Since $|\sin \varphi| \geq \frac{2}{\pi}|\varphi|$, if $-\frac{\pi}{2} \leq \varphi \leq \frac{\pi}{2}$, taking $x, y = 1$, $\phi = \langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle$, $2\pi l = k_0$, we obtain

$$\begin{aligned} |e^{i\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle} - 1| &\geq 2 \left| \sin \frac{\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle - k_0}{2} \right| \\ &\geq \frac{4}{\pi} \left| \frac{\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle - k_0}{2} \right|, \end{aligned}$$

where $k_0 \in \mathbb{Z}$ satisfies $|\frac{\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle - k_0}{2}| \leq \frac{\pi}{2}$.

From Lemma 2.2 we have $|\langle k, \omega_{\nu}(r_{\nu}, \xi) \rangle - k_0| \geq \frac{\delta_{\nu}}{2|k|^{\tau}}$. Also (H3) implies that

$$\begin{aligned} |\langle k, \bar{f}_{\nu}(r_{\nu}, \xi) \rangle| &\leq |k| |\bar{f}_{\nu}(r_{\nu}, \xi)| \leq N_{\nu+1} s_{\nu}^K \delta_{\nu}^K \mu_{\nu} \\ &\leq N_{\nu+1} \frac{\delta_{\nu}}{4N_{\nu+1}^{\tau+1}} \leq \frac{\delta_{\nu}}{4N_{\nu+1}^{\tau}} \leq \frac{\delta_{\nu}}{4|k|^{\tau}}. \end{aligned}$$

Then

$$\begin{aligned} |\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle - k_0| &= |\langle k, \omega_{\nu}(r_{\nu}, \xi) + \bar{f}_{\nu}(r_{\nu}, \xi) \rangle - k_0| \\ &= |\langle k, \omega_{\nu}(r_{\nu}, \xi) \rangle + \langle k, \bar{f}_{\nu}(r_{\nu+1}, \xi) \rangle - k_0| \\ &\geq |\langle k, \omega_{\nu}(r_{\nu}, \xi) \rangle - k_0| - |\langle k, \bar{f}_{\nu}(r_{\nu+1}, \xi) \rangle| \\ &\geq \frac{\delta_{\nu}}{2|k|^{\tau}} - \frac{\delta_{\nu}}{4|k|^{\tau}} \\ &\geq \frac{\delta_{\nu}}{4|k|^{\tau}}, \end{aligned}$$

and therefore,

$$|e^{i\langle k, \omega_{\nu+1}(r_{\nu+1}, \xi) \rangle} - 1| \geq \frac{\delta_{\nu}}{2\pi|k|^{\tau}}.$$

Remark 2.4. Note that $\int_0^{+\infty} x^{\tau+n} e^{-x} dx$ is Γ function and $\Gamma(\tau+n+1) = (\tau+n)!$. Then

$$\begin{aligned} \sum_{0 < |k| \leq N_{\nu+1}} |k|^{\tau} e^{-|k|^{\frac{\rho_{\nu}-\rho_{\nu+1}}{4}}} &\leq \int_0^{N_{\nu+1}} l^{\tau+n} e^{-l^{\frac{\rho_{\nu}-\rho_{\nu+1}}{4}}} dl \\ &\leq \int_0^{+\infty} \left(\frac{4}{\rho_{\nu}-\rho_{\nu+1}}\right)^{\tau+n} x^{\tau+n} e^{-x} d\left(\frac{4}{\rho_{\nu}-\rho_{\nu+1}}x\right) \\ &\leq \frac{4^{\tau+n+1}}{(\rho_{\nu}-\rho_{\nu+1})^{\tau+n+1}} \int_0^{+\infty} x^{\tau+n} e^{-x} dx \\ &= \frac{4^{\tau+n+1}}{(\rho_{\nu}-\rho_{\nu+1})^{\tau+n+1}} \Gamma(\tau+n+1) \\ &= \frac{4^{\tau+n+1}}{(\rho_{\nu}-\rho_{\nu+1})^{\tau+n+1}} (\tau+n)!. \end{aligned}$$

From Cauchy inequality we get

$$\begin{aligned}
& \| \partial^K V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&= \| \partial_{\theta_{\nu+1}}^\kappa \partial_{r_{\nu+1}}^j \partial_\xi^h V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&\leq \cdot \frac{\| V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu)}}{(\frac{\rho_\nu - \rho_{\nu+1}}{4})^\kappa (s_\nu - s_{\nu+1})^j (\sigma_\nu - \sigma_{\nu+1})^h} \\
&\leq \cdot \frac{s_\nu^K \delta_\nu^K \mu_\nu}{(\rho_\nu - \rho_{\nu+1})^{\tau+n+1+\kappa} (\sigma_\nu - \sigma_{\nu+1})^h \delta_\nu s_{\nu+1}^j} \\
&\leq \cdot \frac{s_\nu^{K-j} \delta_\nu^{K-1} \mu_\nu}{(\rho_\nu - \rho_{\nu+1})^{\tau+n+1+\kappa} (\sigma_\nu - \sigma_{\nu+1})^h}. \tag{2.18}
\end{aligned}$$

In an analogous manner we obtain

$$\| U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_3} \leq \cdot \frac{\| f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_4}}{\delta_\nu (\rho_\nu - \rho_{\nu+1})^{\tau+n+1}}, \tag{2.19}$$

and

$$\begin{aligned}
& \| \partial^K U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&= \| \partial_{\theta_{\nu+1}}^\kappa \partial_{r_{\nu+1}}^j \partial_\xi^h U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&\leq \cdot \frac{\| U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu)}}{(\frac{\rho_\nu - \rho_{\nu+1}}{4})^\kappa (s_\nu - s_{\nu+1})^j (\sigma_\nu - \sigma_{\nu+1})^h} \\
&\leq \cdot \frac{s_\nu^K \delta_\nu^K \mu_\nu}{(\rho_\nu - \rho_{\nu+1})^{\tau+n+1+\kappa} (\sigma_\nu - \sigma_{\nu+1})^h \delta_\nu s_{\nu+1}^j} \\
&\leq \cdot \frac{s_\nu^{K-j} \delta_\nu^{K-1} \mu_\nu}{(\rho_\nu - \rho_{\nu+1})^{\tau+n+1+\kappa} (\sigma_\nu - \sigma_{\nu+1})^h}. \tag{2.20}
\end{aligned}$$

Remark 2.5. In order to make sure that

$$\mathcal{U}_{\nu+1} : \mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1}) \subset \mathcal{D}_3 \rightarrow \mathcal{D}_4 \subset \mathcal{D}(\rho_\nu, s_\nu, \sigma_\nu),$$

it requests

$$\begin{aligned}
|r_\nu - r_{\nu+1}| &= \| V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_3} \\
&\leq 4s_{\nu+1} - 3s_{\nu+1} \\
&= s_{\nu+1}, \tag{2.21}
\end{aligned}$$

and

$$|\theta_\nu - \theta_{\nu+1}| = \| U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}_3} \leq \frac{\rho_\nu - \rho_{\nu+1}}{4}. \tag{2.22}$$

That is,

$$(H4) \quad \frac{4^{\tau+n+2} \pi(\tau+n)!}{\delta_\nu (\rho_\nu - \rho_{\nu+1})^{\tau+n+1}} s_\nu^K \delta_\nu^K \mu_\nu \leq s_{\nu+1},$$

and

$$(H5) \quad \frac{4^{\tau+n+2} \pi(\tau+n)!}{\delta_\nu (\rho_\nu - \rho_{\nu+1})^{\tau+n+1}} s_\nu^K \delta_\nu^K \mu_\nu \leq \frac{\rho_\nu - \rho_{\nu+1}}{4}.$$

2.2.3. Estimates of new perturbations

The following lemma shows the estimates of new perturbations in one cycle of the iteration.

Lemma 2.4. *Assume*

$$(H6) \quad 2^{7K} \left(1 + \frac{1}{2^{\nu+1}}\right)^K \mu_\nu^{\frac{2}{3}} \ll 1.$$

Then

$$\|f_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \|g_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \leq s_{\nu+1}^K \delta_{\nu+1}^K \mu_{\nu+1}, \quad (2.23)$$

and

$$\begin{aligned} & \|\partial^K f_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \|\partial^K g_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & \leq \frac{s_{\nu+1}^{K-j} \delta_{\nu+1}^K \mu_{\nu+1}}{(\rho_\nu - \rho_{\nu+1})^\kappa (\sigma_\nu - \sigma_{\nu+1})^h} \end{aligned} \quad (2.24)$$

holds.

Proof. We have already obtained the expressions of $f_{\nu+1}$ and $g_{\nu+1}$ from (2.5) and (2.6), that is,

$$\begin{aligned} & f_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + U_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi) + f_{\nu+1}, r_{\nu+1} + g_{\nu+1}, \xi) \\ & - U_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi), r_{\nu+1} + g_{\nu+1}, \xi) \\ & + U_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi), r_{\nu+1} + g_{\nu+1}, \xi) \\ & - U_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi), r_{\nu+1}, \xi) \\ & = f_\nu(\theta_{\nu+1} + U_{\nu+1}, r_{\nu+1} + V_{\nu+1}, \xi) - f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi) \\ & + R_{N_{\nu+1}} f_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi), \end{aligned} \quad (2.25)$$

$$\begin{aligned} & g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) + V_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi) + f_{\nu+1}, r_{\nu+1} + g_{\nu+1}, \xi) \\ & - V_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi), r_{\nu+1} + g_{\nu+1}, \xi) \\ & + V_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi), r_{\nu+1} + g_{\nu+1}, \xi) \\ & - V_{\nu+1}(\theta_{\nu+1} + \omega_{\nu+1}(r_{\nu+1}, \xi), r_{\nu+1}, \xi) \\ & = g_\nu(\theta_{\nu+1} + U_{\nu+1}, r_{\nu+1} + V_{\nu+1}, \xi) - g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi) + \bar{g}_\nu(r_{\nu+1}, \xi) \\ & + R_{N_{\nu+1}} g_\nu(\theta_{\nu+1}, r_{\nu+1}, \xi). \end{aligned} \quad (2.26)$$

Then $f_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)$ and $g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)$ can be solved by the Implicit Function Theorem. From

$$\begin{aligned} & \|\partial_{\theta_{\nu+1}} U_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \|f_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & + \|\partial_{r_{\nu+1}} U_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \|g_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & + \|f_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & \leq \|\partial_{\theta_{\nu+1}} f_\nu\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \|U_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & + \|\partial_{r_{\nu+1}} f_\nu\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \|V_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & + \|R_{N_{\nu+1}} f_\nu\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})}, \end{aligned}$$

we obtain

$$\begin{aligned}
& \| \partial_{\theta_{\nu+1}} U_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| f_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| \partial_{r_{\nu+1}} U_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| f_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& \leq \cdot \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{\rho_{\nu} - \rho_{\nu+1}} \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{\delta_{\nu}(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} + \cdot \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{s_{\nu} - s_{\nu+1}} \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{\delta_{\nu}(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} \\
& + s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2 \\
& \leq \cdot \frac{s_{\nu}^{2K} \delta_{\nu}^{2K-1} \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+2}} + \cdot \frac{s_{\nu}^{2K-1} \delta_{\nu}^{2K-1} \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} + s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2.
\end{aligned}$$

The intersection property is written as

$$\begin{aligned}
\mathcal{F}_{\nu+1} \mathcal{T}_0 \cap \mathcal{T}_0 &= \mathcal{U}_{\nu+1}^{-1} \circ \mathcal{F}_{\nu} \circ \mathcal{U}_{\nu+1} \mathcal{T}_0 \cap \mathcal{T}_0 \\
&= \mathcal{U}_{\nu+1}^{-1} \circ (\mathcal{F}_{\nu} \circ \mathcal{U}_{\nu+1} \mathcal{T}_0 \cap \mathcal{U}_{\nu+1} \mathcal{T}_0) \\
&\neq \emptyset,
\end{aligned}$$

where $\mathcal{T}_0$ is the initial torus of the mapping. As for $\| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})}$, taking use of the intersection property of twist mapping, there exists a $\theta_{\nu+1}^0$ such that $g_{\nu+1}(\theta_{\nu+1}^0, r_{\nu+1}^0, \xi) = 0$ for each $r_{\nu+1}^0$. Therefore,

$$\begin{aligned}
& \sup \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&= \sup \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) - g_{\nu+1}(\theta_{\nu+1}^0, r_{\nu+1}^0, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&\leq \text{osc}(g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)) \\
&= \text{osc}(g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) - h) \\
&\leq 2 \sup \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) - h \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})},
\end{aligned}$$

where $\text{osc}(g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi))$ denotes the oscillation of function $g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi)$. Specially, taking $h = \bar{g}_{\nu}(r_{\nu+1}, \xi)$, we have

$$\begin{aligned}
& \frac{1}{2} \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
&\leq \sup \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) - \bar{g}_{\nu}(r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})}.
\end{aligned}$$

Therefore, from

$$\begin{aligned}
& \| \partial_{\theta_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| f_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| \partial_{r_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \frac{1}{2} \| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& \leq \| \partial_{\theta_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| f_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| \partial_{r_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) - \bar{g}_{\nu}(r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& \leq \| \partial_{\theta_{\nu+1}} g_{\nu} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| U_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| \partial_{r_{\nu+1}} g_{\nu} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| V_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| R_{N_{\nu+1}} g_{\nu} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})},
\end{aligned}$$

we have

$$\begin{aligned}
& 2 \| \partial_{\theta_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| f_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + 2 \| \partial_{r_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& \leq 2 \| \partial_{\theta_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| f_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + 2 \| \partial_{r_{\nu+1}} V_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + 2 \| g_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi) - \bar{g}_{\nu}(r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& \leq 2 \| \partial_{\theta_{\nu+1}} g_{\nu} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| U_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + 2 \| \partial_{r_{\nu+1}} g_{\nu} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \| V_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + 2 \| R_{N_{\nu+1}} g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi) \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& \leq \cdot \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{\rho_{\nu} - \rho_{\nu+1}} \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{\delta_{\nu}(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} \\
& + \cdot \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{s_{\nu} - s_{\nu+1}} \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}}{\delta_{\nu}(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} \\
& + \cdot s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2 \\
& \leq \cdot \frac{s_{\nu}^{2K} \delta_{\nu}^{2K-1} \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+2}} + \cdot \frac{s_{\nu}^{2K-1} \delta_{\nu}^{2K-1} \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} \\
& + \cdot s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2.
\end{aligned}$$

If

$$(H7) \quad 1 + \frac{3\pi 4^{\tau+n+2}(\tau+n)!}{\delta_{\nu}(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+2}} s_{\nu}^K \delta_{\nu}^K \mu_{\nu} \leq 2,$$

$$(H8) \quad 1 + \frac{3\pi 4^{\tau+n+2}(\tau+n)!}{\delta_{\nu}(s_{\nu} - s_{\nu+1})(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} s_{\nu}^K \delta_{\nu}^K \mu_{\nu} \leq 2,$$

the estimate of new perturbation is

$$\begin{aligned}
& \| f_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \| g_{\nu+1} \|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& \leq \cdot \frac{s_{\nu}^{2K} \delta_{\nu}^{2K-1} \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+2}} + \cdot \frac{s_{\nu}^{2K-1} \delta_{\nu}^{2K-1} \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1}} + \cdot s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2 \\
& \leq \cdot s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2 \\
& \leq \cdot s_{\nu+1}^K \delta_{\nu+1}^K \mu_{\nu+1} \left(\frac{s_{\nu}}{s_{\nu+1}} \right)^K \left(\frac{\delta_{\nu}}{\delta_{\nu+1}} \right)^K \mu_{\nu}^{\frac{2}{3}} \\
& \leq \cdot s_{\nu+1}^K \delta_{\nu+1}^K \mu_{\nu+1} 2^{7K} \left(1 + \frac{1}{2^{\nu+1}} \right)^K \mu_{\nu}^{\frac{2}{3}} \\
& \leq s_{\nu+1}^K \delta_{\nu+1}^K \mu_{\nu+1}.
\end{aligned}$$

Cauchy estimate yields that

$$\begin{aligned}
& \| \partial^K f_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \| \partial^K g_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& = \| \partial_{\theta_{\nu+1}}^{\kappa} \partial_{r_{\nu+1}}^j \partial_{\xi}^h f_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\
& + \| \partial_{\theta_{\nu+1}}^{\kappa} \partial_{r_{\nu+1}}^j \partial_{\xi}^h g_{\nu+1} \|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})}
\end{aligned}$$

$$\leq \frac{\|f_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \|g_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})}}{(\rho_{\nu} - \rho_{\nu+1})^{\kappa} (s_{\nu} - s_{\nu+1})^j (\sigma_{\nu} - \sigma_{\nu+1})^h},$$

that is,

$$\begin{aligned} & \|\partial^K f_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \|\partial^K g_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & \leq \frac{s_{\nu}^K \delta_{\nu}^K \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\kappa} s_{\nu+1}^j (\sigma_{\nu} - \sigma_{\nu+1})^h} \\ & \leq \frac{s_{\nu}^{K-j} \delta_{\nu}^K \mu_{\nu}^2}{(\rho_{\nu} - \rho_{\nu+1})^{\kappa} (\sigma_{\nu} - \sigma_{\nu+1})^h} \\ & \leq \frac{s_{\nu+1}^{K-j} \delta_{\nu+1}^K \mu_{\nu+1}}{(\rho_{\nu} - \rho_{\nu+1})^{\kappa} (\sigma_{\nu} - \sigma_{\nu+1})^h} \left(\frac{s_{\nu}}{s_{\nu+1}}\right)^{K-j} \left(\frac{\delta_{\nu}}{\delta_{\nu+1}}\right)^K \mu_{\nu}^{\frac{2}{3}} \\ & \leq \frac{s_{\nu+1}^{K-j} \delta_{\nu+1}^K \mu_{\nu+1}}{(\rho_{\nu} - \rho_{\nu+1})^{\kappa} (\sigma_{\nu} - \sigma_{\nu+1})^h} 2^{7K-7j} \left(1 + \frac{1}{2^{\nu+1}}\right)^K \mu_{\nu}^{\frac{2}{3}} \\ & \leq \frac{s_{\nu+1}^{K-j} \delta_{\nu+1}^K \mu_{\nu+1}}{(\rho_{\nu} - \rho_{\nu+1})^{\kappa} (\sigma_{\nu} - \sigma_{\nu+1})^h}. \end{aligned}$$

□

Remark 2.6. It allows us to consider the problem in the whole domain in view of Whitney smoothness in a standard manner, and thus we omit the details.

3. Proof of Theorem 1.1

In this subsection, Iteration lemma is used to prove Theorem 1.1. Let $s_0, \delta_0, \mu_0, \rho_0, \sigma_0, \epsilon_0, f_0, g_0, \mathcal{F}_0, \mathcal{D}(\rho_0, s_0, \sigma_0)$ be given in subsection 2.1, we define

$$\begin{aligned} \epsilon_{\nu} &= s_{\nu}^K \delta_{\nu}^K \mu_{\nu}, & s_{\nu} &= \frac{1}{2^{7\nu+2}} s_0, & \delta_{\nu} &= \left(\frac{1}{2} + \frac{1}{2^{\nu+1}}\right) \delta_0, & \mu_{\nu} &= \mu_{\nu-1}^{\frac{4}{3}}, \\ \rho_{\nu} &= \left(\frac{1}{2} + \frac{1}{2^{\nu+1}}\right) \rho_0, & \sigma_{\nu} &= \frac{1}{2^{7\nu+2}} \sigma_0, & N_{\nu+1} &= ([\log \frac{1}{\mu_{\nu}}] + 1)^{3\chi}, \\ \mathcal{D}_j &= \mathcal{D}(\rho_{\nu+1} + \frac{j-1}{4}(\rho_{\nu} - \rho_{\nu+1}), j s_{\nu+1}, j \sigma_{\nu+1}), & j &= 1, 2, 3, 4, \\ \hat{\mathcal{D}}(\rho_{\nu}, s_{\nu}, \sigma_{\nu}) &= \mathcal{D}(\rho_{\nu+1} + \frac{3}{4}(\rho_{\nu} - \rho_{\nu+1}), s_{\nu+1}, \sigma_{\nu+1}), & \mathcal{D}_1 &= \mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1}), \end{aligned}$$

for any $\nu = 1, 2, \dots$.

The Iteration lemma is stated as follows.

3.1. Iteration Lemma

Lemma 3.1. Assume (H1)-(H8). Then for mapping $\mathcal{F}_{\nu}$ on $\mathcal{D}(\rho_{\nu}, s_{\nu}, \sigma_{\nu})$, if

$$\|f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}(\rho_{\nu}, s_{\nu}, \sigma_{\nu})} + \|g_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\mathcal{D}(\rho_{\nu}, s_{\nu}, \sigma_{\nu})} \leq s_{\nu}^K \delta_{\nu}^K \mu_{\nu},$$

the following hold for all $\nu = 1, 2, \dots$.

(i) There exists a transformation $\mathcal{U}_{\nu+1}$:

$$\begin{cases} \theta_{\nu} = \theta_{\nu+1} + U_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi), \\ r_{\nu} = r_{\nu+1} + V_{\nu+1}(\theta_{\nu+1}, r_{\nu+1}, \xi), \end{cases}$$

such that

$$\mathcal{U}_{\nu+1} \circ \mathcal{F}_{\nu+1} = \mathcal{F}_{\nu} \circ \mathcal{U}_{\nu+1},$$

and

$$\|\partial^K U_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \leq \frac{s_{\nu}^{K-j} \delta_{\nu}^{K-1} \mu_{\nu}}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1+\kappa} (\sigma_{\nu} - \sigma_{\nu+1})^h}, \quad (3.1)$$

$$\|\partial^K V_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \leq \frac{s_{\nu}^{K-j} \delta_{\nu}^{K-1} \mu_{\nu}}{(\rho_{\nu} - \rho_{\nu+1})^{\tau+n+1+\kappa} (\sigma_{\nu} - \sigma_{\nu+1})^h}; \quad (3.2)$$

(ii)

$$\|f_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \|g_{\nu+1}\|_{\mathcal{D}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \leq s_{\nu+1}^K \delta_{\nu+1}^K \mu_{\nu+1}, \quad (3.3)$$

$$\begin{aligned} & \|\partial^K f_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} + \|\partial^K g_{\nu+1}\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & \leq \frac{s_{\nu+1}^{K-j} \delta_{\nu+1}^K \mu_{\nu+1}}{(\rho_{\nu} - \rho_{\nu+1})^{\kappa} (\sigma_{\nu} - \sigma_{\nu+1})^h}; \end{aligned} \quad (3.4)$$

(iii)

$$|\partial_{r_{\nu+1}}^j \partial_{\xi}^h (\omega_{\nu+1}(r_{\nu+1}, \xi) - \omega_{\nu}(r_{\nu}, \xi))| \leq \frac{s_{\nu}^{K-j} \delta_{\nu}^K \mu_{\nu}}{(\sigma_{\nu} - \sigma_{\nu+1})^h}. \quad (3.5)$$

Proof. It is easy to see that we can take sufficiently small ϵ such that (H1)-(H8) hold. Now (i) follows from Lemma 2.3, (2.18), and (2.20), (ii) follows from Lemma 2.4. As for (iii), we have

$$\begin{aligned} & |\partial_{r_{\nu+1}}^j \partial_{\xi}^h (\omega_{\nu+1}(r_{\nu+1}, \xi) - \omega_{\nu}(r_{\nu}, \xi))| \\ & \leq \|\partial_{r_{\nu+1}}^j \partial_{\xi}^h \bar{f}_{\nu}(r_{\nu+1}, \xi)\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & = \frac{1}{(2\pi)^n} \|\partial_{r_{\nu+1}}^j \partial_{\xi}^h \left(\int_{\mathbb{T}^n} f_{\nu}(\theta_{\nu+1}, r_{\nu+1}) d\theta_{\nu+1} \right)\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} \\ & \leq \frac{1}{(2\pi)^n} \int_{\mathbb{T}^n} \|\partial_{r_{\nu+1}}^j \partial_{\xi}^h f_{\nu}(\theta_{\nu+1}, r_{\nu+1}, \xi)\|_{\hat{\mathcal{D}}(\rho_{\nu+1}, s_{\nu+1}, \sigma_{\nu+1})} d\theta_{\nu+1} \\ & \leq \frac{s_{\nu}^{K-j} \delta_{\nu}^K \mu_{\nu}}{(\sigma_{\nu} - \sigma_{\nu+1})^h}. \end{aligned}$$

Above all, the proof is complete. $\square$

3.2. Convergence

Denote $\mathcal{I}_{\nu}$ by

$$\mathcal{I}_{\nu} = \mathcal{U}_1 \circ \mathcal{U}_2 \circ \cdots \circ \mathcal{U}_{\nu},$$

then

$$\mathcal{F}_{\nu} = \mathcal{I}_{\nu}^{-1} \circ \mathcal{F}_0 \circ \mathcal{I}_{\nu}.$$

Note that $\mathcal{I}_{\nu}$ has the form

$$\begin{cases} \theta_0 = \theta_{\nu} + P_{\nu}(\theta_{\nu}, r_{\nu}, \xi) = \theta_{\nu} + \sum_{l=1}^{\nu} U_l(\theta_l, r_l, \xi), \\ r_0 = r_{\nu} + Q_{\nu}(\theta_{\nu}, r_{\nu}, \xi) = r_{\nu} + \sum_{l=1}^{\nu} V_l(\theta_l, r_l, \xi), \end{cases}$$

and

$$\begin{aligned}
& \|P_l(\theta_l, r_l, \xi)\|_{\mathcal{D}(\rho_l, s_l, \sigma_l)} + \|Q_l(\theta_l, r_l, \xi)\|_{\mathcal{D}(\rho_l, s_l, \sigma_l)} \\
& \leq \sum_{l=1}^{\nu} (\|U_l(\theta_l, r_l, \xi)\|_{\mathcal{D}(\rho_l, s_l, \sigma_l)} + \|V_l(\theta_l, r_l, \xi)\|_{\mathcal{D}(\rho_l, s_l, \sigma_l)}) \\
& \leq \sum_{l=1}^{\nu} \frac{\|f_{l-1}(\theta_l, r_l, \xi)\|_{\mathcal{D}(\rho_l, s_l, \sigma_l)} + \|g_{l-1}(\theta_l, r_l, \xi)\|_{\mathcal{D}(\rho_l, s_l, \sigma_l)}}{\delta_{l-1}(\rho_{l-1} - \rho_l)^{\tau+n+1}} \\
& \leq \sum_{l=1}^{\nu} \frac{s_{l-1}^K \delta_{l-1}^{K-1} \mu_{l-1}}{(\rho_{l-1} - \rho_l)^{\tau+n+1}} \\
& \leq \frac{s_0^K \delta_0^{K-1} \mu_0}{\rho_0^{\tau+n+1}}.
\end{aligned}$$

Therefore, P_l and Q_l are uniformly convergent on $\mathcal{D}(\rho_\infty, s_\infty, \sigma_\infty)$. Let $P = \lim_{l \rightarrow \infty} P_l$, $Q = \lim_{l \rightarrow \infty} Q_l$, then $\mathcal{I}_\infty$ becomes

$$\begin{cases} \theta_0 = \theta_\infty + P(\theta, r, \xi), \\ r_0 = r_\infty + Q(\theta, r, \xi). \end{cases}$$

Moreover, from

$$\omega_{\nu+1}(r_{\nu+1}, \xi) = \omega_\nu(r_\nu, \xi) + \bar{f}_\nu(r_{\nu+1}, \xi),$$

and

$$|\partial_{r_{\nu+1}}^j \partial_\xi^h (\omega_{\nu+1}(r_{\nu+1}, \xi) - \omega_\nu(r_\nu, \xi))| \leq \frac{s_\nu^{K-j} \delta_\nu^K \mu_\nu}{(\sigma_\nu - \sigma_{\nu+1})^h},$$

we get

$$\begin{aligned}
|\partial_{r_{\nu+1}}^j \partial_\xi^h (\omega_\infty(r_\infty, \xi) - \omega_0(r_0, \xi))| & \leq |\partial_{r_{\nu+1}}^j \partial_\xi^h (\sum_{\nu=0}^{\infty} (\omega_{\nu+1}(r_{\nu+1}, \xi) - \omega_\nu(r_\nu, \xi)))| \\
& \leq \sum_{\nu=0}^{\infty} |\partial_{r_{\nu+1}}^j \partial_\xi^h (\omega_{\nu+1}(r_{\nu+1}, \xi) - \omega_\nu(r_\nu, \xi))| \\
& \leq \sum_{\nu=0}^{\infty} \frac{s_\nu^{K-j} \delta_\nu^K \mu_\nu}{(\sigma_\nu - \sigma_{\nu+1})^h} \\
& \leq \frac{s_0^{K-j} \delta_0^K \mu_0}{\sigma_0^h},
\end{aligned}$$

that is, ω_ν is uniformly convergent to $\omega_\infty(r_\infty, \xi)$ on $\mathcal{D}(\rho_\infty, s_\infty, \sigma_\infty)$. Therefore $\mathcal{I}_\infty$ transforms $\mathcal{F}_0$ into

$$\begin{cases} \theta_\infty^1 = \theta_\infty + \omega_\infty(r_\infty, \xi), \\ r_\infty^1 = r_\infty. \end{cases}$$

3.3. Measure estimates

The KAM steps are independent of the relationship between n and $m + p$, but measure estimates are not. For this purpose, let us consider the following three cases.

Case 1: $n = m + p$.

We assume $\tilde{\omega} = (\omega, -1)$, $\tilde{r} = (r, \xi, r_{n+1})$, $\tilde{k} = (k, k_0)$, $\tilde{D} = D \times G \times (1, 2)$, and

$$\tilde{D}_\delta = \{\tilde{r} \in \tilde{D} : |\langle \tilde{k}, \tilde{\omega} \rangle| \geq \delta |k|^{-\tau}, \forall \tilde{k} \in \mathbb{Z}^{n+1} \setminus \{0\}\},$$

then $\text{meas}((D \times G) \setminus D_\delta)$ is equivalent to $\text{meas}(\tilde{D} \setminus \tilde{D}_\delta)$.

Clearly on $\tilde{D}$,

$$\text{rank}\{\partial_{\tilde{r}}^\alpha \tilde{\omega}(\tilde{r}) : 0 \leq |\alpha| \leq K-1\} = n+1,$$

we have

$$\text{meas}(\tilde{D} \setminus \tilde{D}_\delta) \leq c_0 \delta^{\frac{1}{K}},$$

that is

$$\text{meas}((D \times G) \setminus D_\delta) \leq c_0 \delta^{\frac{1}{K}},$$

where c_0 is a constant. See Appendix for details.

Case 2: $m + p < n$.

Assume $\tilde{r} = (r, \xi, r_{m+p+1}, \dots, r_n, r_{n+1})$, $\tilde{\omega}(\tilde{r}) = (\omega, -1)$, $\tilde{k} = (k, k_0)$, then

$$\text{rank}\{\partial_{\tilde{r}}^\alpha \tilde{\omega}(\tilde{r}) : 0 \leq |\alpha| \leq K-1\} = n+1,$$

and $\tilde{D} = D \times G \times \underbrace{(1, 2) \times \dots \times (1, 2)}_{n+1-m-p}$, $\tilde{D}_\delta = D_\delta \times \underbrace{(1, 2) \times \dots \times (1, 2)}_{n+1-m-p}$. Then

$$|\langle \tilde{k}, \tilde{\omega} \rangle| \geq \delta |k|^{-\tau}, \quad \forall \tilde{k} \in \mathbb{Z}^{n+1} \setminus \{0\},$$

on $\tilde{D}_\delta$ is equivalent to the Diophantine condition

$$|\langle k, \omega \rangle - k_0| \geq \delta |k|^{-\tau}, \quad \forall k \in \mathbb{Z}^n \setminus \{0\}, \quad \forall k_0 \in \mathbb{Z},$$

on D_δ .

Then from the estimates of measure, we obtain

$$\text{meas}(D \setminus D_\delta) \leq c_0 \delta^{\frac{1}{K}}.$$

Case 3: $n < m + p$.

We have known that

$$D_\delta = \{(r, \xi) \in D \times G : |\langle k, \omega(r, \xi) \rangle - k_0| \geq \frac{\delta}{|k|^\tau}, \forall k \in \mathbb{Z}^n \setminus \{0\}, \forall k_0 \in \mathbb{Z}\}.$$

Assume

$$D_\delta^1 = \{(r, \xi) \in D \times G : |\langle k, \omega \rangle + \langle \check{k}, \check{r} \rangle - k_0| \geq \frac{\delta}{|k|^\tau}, \\ \forall k \in \mathbb{Z}^n \setminus \{0\}, \forall \check{k} \in \mathbb{Z}^{m+p-n}, \forall k_0 \in \mathbb{Z}\},$$

where $\check{k} = (k_{n+1}, k_{n+2}, \dots, k_m, \dots, k_{m+p})$, $\check{r} = (r_{n+1}, r_{n+2}, \dots, r_m, \xi)$, if $n < m$, (resp. $\check{k} = (k_{n+1}, \dots, k_{m+p})$, $\check{r} = (\xi_{n-m+1}, \dots, \xi_p)$, if $m < n < m + p$). Obviously, $D_\delta^1 \subset D_\delta$, therefore,

$$\text{meas}((D \times G) \setminus D_\delta) \leq \text{meas}((D \times G) \setminus D_\delta^1).$$

From the non-degenerate condition (A2), we have

$$\text{rank}\left\{\frac{\partial^\alpha \omega}{\partial(r, \xi)^\alpha} : 0 \leq |\alpha| \leq K-1\right\} = n.$$

Let $\omega' = (\omega, r_{n+1}, r_{n+2}, \dots, r_m, \xi)$, (resp. $\omega' = (\omega, \xi_{n-m+1}, \dots, \xi_p)$), it is easy to see that

$$\text{rank}\left\{\frac{\partial^\alpha \omega'}{\partial(r, \xi)^\alpha} : 0 \leq |\alpha| \leq K-1\right\} = m+p.$$

If $\tilde{\omega} = (\omega', -1)$, $\tilde{r} = (r, \xi, r_{m+p+1})$, $\tilde{k} = (k, \tilde{k}, k_0)$, $\tilde{D} = D \times G \times (1, 2)$ and $\tilde{D}_\delta = D_\delta^1 \times (1, 2)$, that is,

$$\tilde{D}_\delta = \{\tilde{r} \in \tilde{D} : |\langle \tilde{k}, \tilde{\omega} \rangle| \geq \frac{\delta}{|\tilde{k}|^\tau}, \forall \tilde{k} \in \mathbb{Z}^{m+p+1} \setminus \{0\}\},$$

one has

$$\text{rank}\{\partial_{\tilde{r}}^\alpha \tilde{\omega}(\tilde{r}) : 0 \leq |\alpha| \leq K-1\} = m+p+1.$$

Hence, from the Appendix we get

$$\text{meas}(\tilde{D} \setminus \tilde{D}_\delta) \leq c_0 \delta^{\frac{1}{K}},$$

and

$$\text{meas}((D \times G) \setminus D_\delta^1) = \text{meas}(\tilde{D} \setminus \tilde{D}_\delta).$$

Therefore,

$$\text{meas}((D \times G) \setminus D_\delta) \leq \text{meas}((D \times G) \setminus D_\delta^1) = \text{meas}(\tilde{D} \setminus \tilde{D}_\delta) \leq c_0 \delta^{\frac{1}{K}}.$$

4. Appendix

Here we use the method of measure estimates mentioned in [15, p1787].

Assume $\tilde{\omega} = (\omega, -1)$, $\tilde{r} = (r, r_{n+1})$, $\tilde{k} = (k, k_0)$, $\tilde{D} = D \times (1, 2)$, and

$$\tilde{D}_\delta = D_\delta \times (1, 2) = \{\tilde{r} \in \tilde{D} : |\langle \tilde{k}, \tilde{\omega} \rangle| \geq \delta |k|^{-\tau}, \forall \tilde{k} \in \mathbb{Z}^{n+1} \setminus \{0\}\}.$$

Then $\text{meas}(D \setminus D_\delta)$ is equal to $\text{meas}(\tilde{D} \setminus \tilde{D}_\delta)$.

Assume

$$\mathcal{R} = \{\tilde{r} \in \tilde{D} : |\langle \tilde{k}, \tilde{\omega} \rangle| \leq \delta |k|^{-\tau}, \forall \tilde{k} \in \mathbb{Z}^{n+1} \setminus \{0\}\},$$

and

$$z(\tilde{r}) = \langle \zeta, \tilde{\omega} \rangle, \quad \zeta = \frac{\tilde{k}}{|\tilde{k}|} \in S^{n+1},$$

where S^{n+1} is a $(n+1)$ -dimensional unit sphere. The Taylor expansion of $z(\tilde{r})$ is

$$z(\tilde{r}) = \zeta^T \Lambda(\tilde{r}) \tilde{r},$$

where,

$$\Lambda(\tilde{r}) = \left(\tilde{\omega}(\tilde{r}), \partial_{\tilde{r}} \tilde{\omega}(\tilde{r}), \dots, \partial_{\tilde{r}}^n \tilde{\omega}(\tilde{r}), \int_0^1 (1-t)^{|n+1|} \partial_{\tilde{r}}^{n+1} \tilde{\omega}(\tilde{r}_0 + (\tilde{r}_0 + t\tilde{r}_0)) dt \right),$$

$$\tilde{r} = (1, \hat{r}, \dots, \hat{r}^n, \hat{r}^{n+1}),$$

$$\hat{r} = \tilde{r} - \tilde{r}_0.$$

By the non-degenerate condition (A2), $\text{rank}(\Lambda(\tilde{r})) = n + 1$ for $\tilde{r} \in \tilde{D}$ when $|\alpha| = K$. Then there exists an orthogonal matrix $O_{\tilde{r}_0} = (o_{ij})$ whose elements in different columns and rows are only one is 1 and others are 0, such that

$$\Lambda(\tilde{r})O_{\tilde{r}_0} = (A(\tilde{r}), B(\tilde{r})),$$

where $\det A(\tilde{r}) \neq 0$ and $\tilde{r}$ is in the neighborhood $\bar{D} \subset \mathbb{R}^n$ of $\tilde{r}_0$. Therefore,

$$z(\tilde{r}) = \zeta^T \Lambda(\tilde{r}) \bar{r} = \zeta^T \Lambda(\tilde{r}) O_{\tilde{r}_0}^{-1} \bar{r} = (\zeta^T A(\tilde{r}), \zeta^T B(\tilde{r})) O_{\tilde{r}_0}^{-1} \bar{r},$$

and

$$z^T(\tilde{r}) = \bar{r}^T O_{\tilde{r}_0} (\zeta^T \Lambda(\tilde{r}) O_{\tilde{r}_0})^T = \bar{r}^T O_{\tilde{r}_0} \begin{pmatrix} A(\tilde{r}) \zeta \\ B(\tilde{r}) \zeta \end{pmatrix}.$$

From the non-degenerate condition

$$\text{rank}\{\partial_{\tilde{r}}^\alpha z(\tilde{r}) : 0 \leq |\alpha| \leq K - 1\} = 1,$$

which is equivalent to the non-degenerate condition that appears in this paper, we get

$$\text{rank} \begin{pmatrix} A^T(\tilde{r}) \zeta \zeta^T A(\tilde{r}) & A^T(\tilde{r}) \zeta \zeta^T B(\tilde{r}) \\ B^T(\tilde{r}) \zeta \zeta^T A(\tilde{r}) & B^T(\tilde{r}) \zeta \zeta^T B(\tilde{r}) \end{pmatrix} = 1, \quad \tilde{r} \in \bar{D}_{\tilde{r}_0}, \quad \zeta \in S^{n+1}.$$

Therefore, for all $\tilde{r}$ in the neighborhood of $\tilde{r}_0$, there exists an orthogonal matrix $O_{\tilde{r}}$ continuously depending on $\tilde{r}$ such that

$$O_{\tilde{r}}^{-1} \begin{pmatrix} A^T(\tilde{r}) \zeta \zeta^T A(\tilde{r}) & A^T(\tilde{r}) \zeta \zeta^T B(\tilde{r}) \\ B^T(\tilde{r}) \zeta \zeta^T A(\tilde{r}) & B^T(\tilde{r}) \zeta \zeta^T B(\tilde{r}) \end{pmatrix} O_{\tilde{r}} = \text{diag}(0, \dots, 0, \underbrace{\lambda(\tilde{r}, \zeta)}_{i\text{-th}}, 0, \dots, 0),$$

where the position of λ depends on ζ , and

$$|(O_{\tilde{r}_0}^{-1} \bar{r})_i| = |(O_{\tilde{r}_0})_i \bar{r}| \geq \min_j |\hat{r}_j|^K.$$

Taking use of the Poincaré Separation Theorem, we have

$$\begin{aligned} z^T(\tilde{r}) z(\tilde{r}) &= \bar{r}^T O_{\tilde{r}_0} O_{\tilde{r}_0}^{-1} \Lambda^T(\tilde{r}) \zeta \zeta^T \Lambda(\tilde{r}) O_{\tilde{r}_0} O_{\tilde{r}_0}^{-1} \bar{r} \\ &= \bar{r}^T O_{\tilde{r}_0} \begin{pmatrix} A^T(\tilde{r}) \zeta \zeta^T A(\tilde{r}) & A^T(\tilde{r}) \zeta \zeta^T B(\tilde{r}) \\ B^T(\tilde{r}) \zeta \zeta^T A(\tilde{r}) & B^T(\tilde{r}) \zeta \zeta^T B(\tilde{r}) \end{pmatrix} O_{\tilde{r}_0}^{-1} \bar{r} \\ &\geq \bar{r}^T O_{\tilde{r}_0} \text{diag}(0, \dots, 0, \lambda(\tilde{r}, \zeta), 0, \dots, 0) O_{\tilde{r}_0}^{-1} \bar{r} \\ &= \min \lambda(\tilde{r}, \zeta) |(O_{\tilde{r}_0} \bar{r})_i|^2 \\ &\geq \lambda_{\tilde{r}_0} (\min_j |\hat{r}_j|)^{2K}. \end{aligned}$$

Then

$$\text{meas}(\{\tilde{r} \in \tilde{D} \cap \bar{D}_{\tilde{r}_0} \mid |\langle \zeta, \tilde{\omega} \rangle|^2 \leq \delta^2 |k|^{-2\tau}\})$$

$$\begin{aligned}
&< \text{meas}(\{\tilde{r} \in \tilde{D} \cap \bar{D}_{\tilde{r}_0} \mid \lambda_{\tilde{r}_0}(\min_j |\hat{r}_j|)^{2K+2} \leq \delta^2 |k|^{-2\tau}\}) \\
&\leq \frac{1}{\lambda_{\tilde{r}_0}} (\text{diam} \tilde{D})^n \delta^{\frac{1}{K}} |k|^{-\frac{\tau}{K}} \\
&\leq c_0 \delta^{\frac{1}{K}},
\end{aligned}$$

where $\text{diam} \tilde{D}$ is the diameter of area $\tilde{D}$ and c_0 is a constant depends on $\tilde{D}$, $\text{diam} \tilde{D}$, n , and $\lambda_{\tilde{r}_0}$, that is,

$$\text{meas}(\mathcal{R}) \leq c_0 \delta^{\frac{1}{K}}.$$

Thus

$$\text{meas}(\tilde{D} \setminus \tilde{D}_\delta) \leq \text{meas}(\mathcal{R}) \leq c_0 \delta^{\frac{1}{K}},$$

due to $\tilde{D} \setminus \tilde{D}_\delta \subset \mathcal{R}$, and

$$\text{meas}(D \setminus D_\delta) \leq c_0 \delta^{\frac{1}{K}}.$$

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References

- [1] V. I. Arnol'd, *Proof of A. N. Kolmogorov's theorem on the preservation of quasi-periodic motions under small perturbations of the Hamiltonian*, Russian Math. Surveys, 1963, 18, 9–36.
- [2] C. Cheng and Y. Sun, *Existence of invariant tori in three-dimensional measure-preserving mappings*, Celestial Mech. Dynam. Astronom., 1989/90, 47(3), 275–292.
- [3] F. Cong, Y. Li and M. Huang, *Invariant tori for nearly twist mappings with intersection property*, Northeast. Math. J., 1996, 12(3), 280–298.
- [4] L. Chierchia and C. E. Koudjina, *V. I. Arnold's "Global" KAM theorem and geometric measure estimates*, Regul. Chaotic Dyn., 2021, 26(1), 61–88.
- [5] R. C. Calleja, A. Celletti and R. de la Llave, *KAM quasi-periodic solutions for the dissipative standard map*, Commun. Nonlinear Sci. Numer. Simul., 2022, 106(106111), 1–29.
- [6] M. R. Herman, *Sur les courbes invariantes par les difféomorphismes de l'anneau*, Astérisque, 1986, 2(144), 1–248.
- [7] A. N. Kolmogorov, *On conservation of conditionally periodic motions for a small change in Hamilton's function*, Dokl. Akad. Nauk SSSR, 1954, 98, 527–530.
- [8] C. E. Koudjina, *A KAM theorem for finitely differentiable Hamiltonian systems*, J. Differential Equations, 2020, 269(6), 4720–4750.
- [9] M. Levi and J. Moser, *A Lagrangian proof of the invariant curve theorem for twist mapping*, Smooth ergodic theory and its applications. Proc. Sympos. Pure Math. Amer. Math. Soc., Providence, RI, 2001, 69, 733–746.

- [10] Y. Li and Y. Yi, *Persistence of invariant tori in generalized Hamiltonian systems*, Ergodic Theory Dynam. Systems, 2002, 22(4), 1233–1261.
- [11] J. Moser, *On invariant curves of area-preserving mappings of an annulus*, Nachr. Akad. Wiss. Göttingen Math. Phys. Kl. II, 1962 (1962), 1–20.
- [12] J. Moser, *A rapidly convergent iteration method and non-linear partial differential equations. I*, Ann. Scuola Norm. Sup. Pisa Cl. Sci., 1966, 20(3), 265–315.
- [13] J. Moser, *A rapidly convergent iteration method and non-linear partial differential equations. II*, Ann. Scuola Norm. Sup. Pisa Cl. Sci., 1966, 20(3), 499–535.
- [14] J. Moser, *On the construction of almost periodic solutions for ordinary differential equations*, 1970 Proc. Internat. Conf. on Functional Analysis and Related Topics, 1969, 60–67.
- [15] W. Qian, Y. Li and X. Yang, *Melnikov's conditions in matrices*, J. Dynam. Differential Equations, 2020, 32(4), 1779–1795.
- [16] H. Rüssmann, *On the existence of invariant curves of twist mappings of an annulus*, Geometric dynamics, 677–718, Lecture Notes in Math., 1007, Springer, Berlin, 1983.
- [17] N. V. Svanidze, *Small perturbations of an integrable dynamical system with an integral invariant*, Trudy Mat. Inst. Steklov., 1980, 147, 124–146.
- [18] M. B. Sevryuk, *Partial preservation of the frequencies and Floquet exponents of invariant tori in KAM theory reversible context 2*, J. Math. Sci. (N. Y.), 2021, 253(5), 730–753.
- [19] F. Trujillo, *Uniqueness properties of the KAM curve*, Discrete Contin. Dyn. Syst., 2021, 41(11), 5165–5182.
- [20] Z. Xia, *Existence of invariant tori in volume-preserving diffeomorphisms*, Ergodic Theory Dynam. Systems, 1992, 12(3), 621–631.
- [21] L. Yang and X. Li, *Existence of periodically invariant tori on resonant surfaces for twist mappings*, Discrete Contin. Dyn. Syst., 2020, 40(3), 1389–1409.
- [22] E. Zehnder, *Generalized implicit function theorems with applications to some small divisor problems. I*, Comm. Pure Appl. Math., 1975, 28, 91–140.
- [23] E. Zehnder, *Generalized implicit function theorems with applications to some small divisor problems. II*, Comm. Pure Appl. Math., 1976, 29(1), 49–111.